

1965 Nobel Prize in Physics: Tomonaga, Schwinger, Feynman

Chirag Gokani

PHYS 4352 - Concepts of Modern Physics

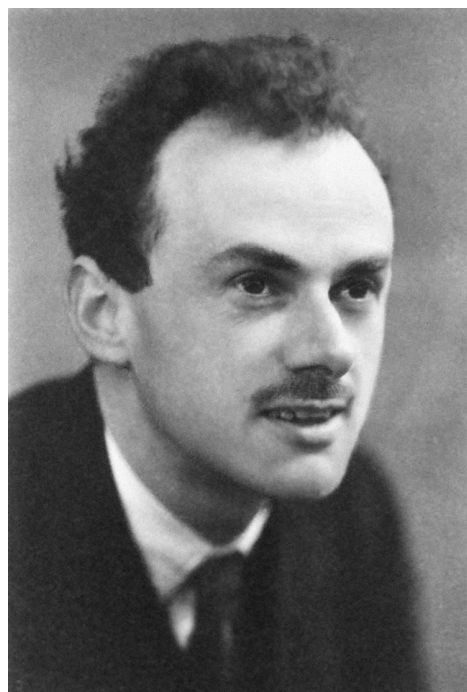
UNIVERSITY OF TEXAS AT DALLAS

November 17, 2020

...for their fundamental work in quantum electrodynamics, with deep-ploughing consequences for the physics of elementary particles.



$$(\square + \mu^2)\psi = 0$$

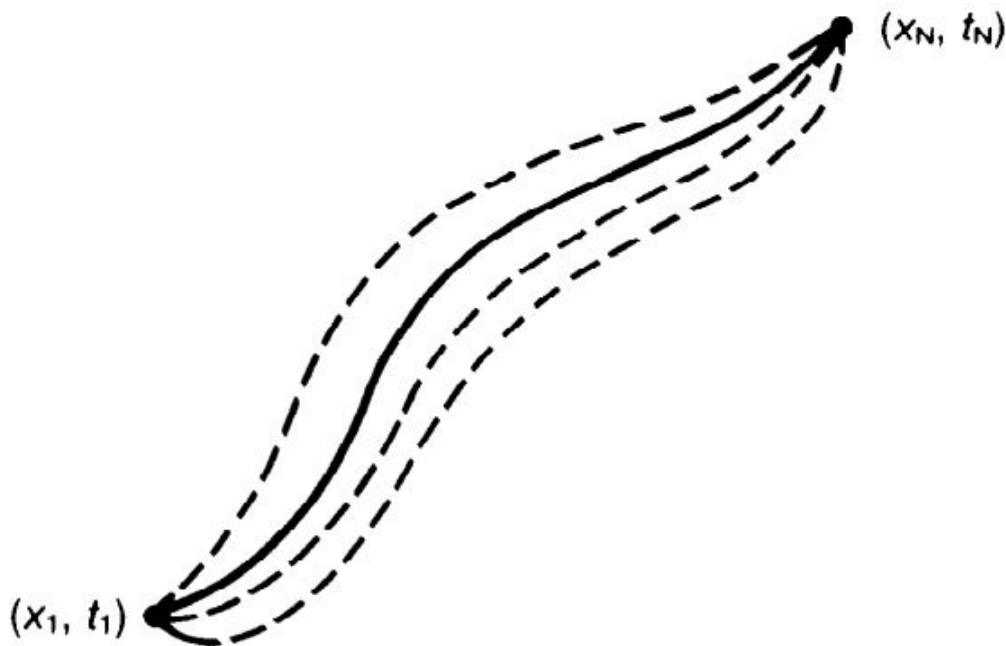


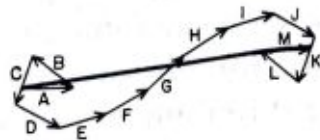
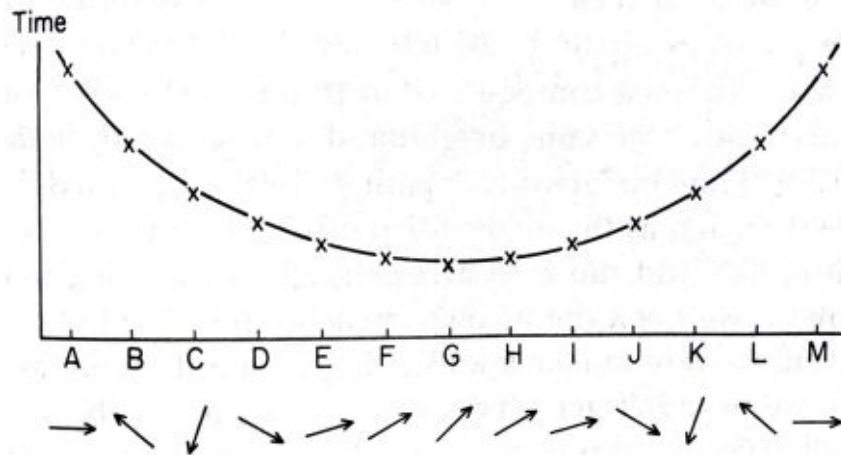
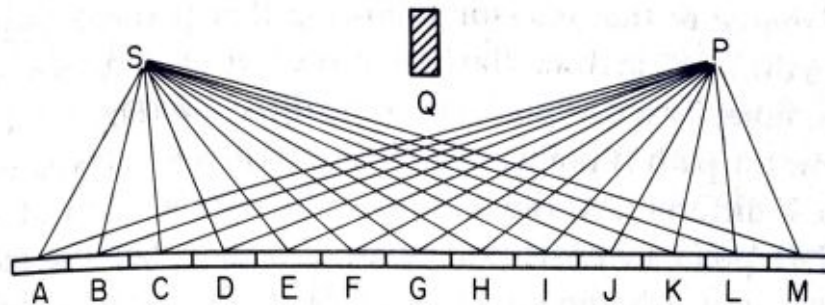
1.00118

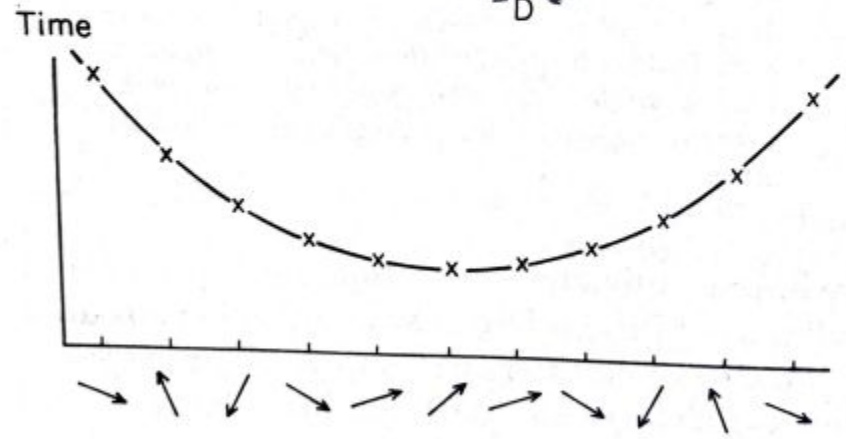
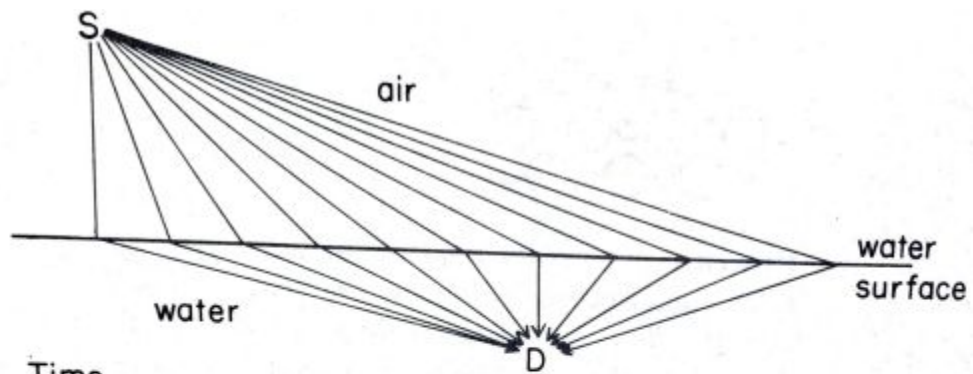
∞

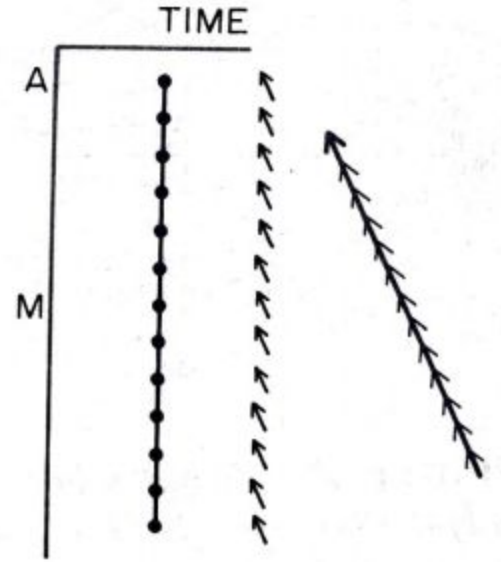
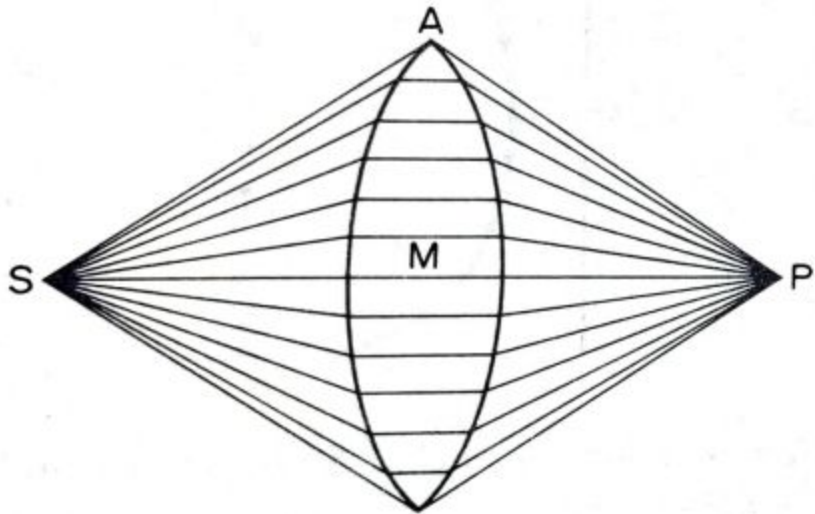
Schrödinger...
$$i\hbar \frac{\partial}{\partial t} \Psi(\vec{x}', t) = - \left(\frac{\hbar^2}{2m} \right) \nabla'^2 \Psi(\vec{x}', t) + V(\vec{x}') \Psi(\vec{x}', t)$$

Heisenberg...
$$\frac{dA(t)}{dt} = \frac{i}{\hbar} [H, A(t)] + \frac{\partial A}{\partial t}$$









$$\exp \left(i \int_{t_1}^{t_2} \frac{dt L_{\text{classical}}(x, \dot{x})}{\hbar} \right) \text{ corresponds to } \langle x_2, t_2 | x_1, t_1 \rangle$$

Calling $S(n, n-1) \equiv \int_{t_{n-1}}^{t_n} dt L_{\text{classical}}(x, \dot{x})$, we can apply successive Dirac-like expressions that somehow contributes to $\langle x_N, t_N | x_1, t_1 \rangle$:

$$\prod_{n=2}^N \exp \left(\frac{iS(n, n-1)}{\hbar} \right) = \dots = \exp \left(\frac{iS(N, 1)}{\hbar} \right)$$

Feynman noted that the classical path is recovered in the $\hbar \rightarrow 0$ limit, leading him to conjecture

$$\langle x_N, t_N | x_1, t_1 \rangle \sim \sum_{\text{all paths } (A, B, \dots)} \exp \left(\frac{iS(N, 1)}{\hbar} \right)$$

Inserting a weighting factor $\frac{1}{w(\Delta t)}$ and considering the case that the time interval $t_n - t_{n-1}$ is infinitesimally small,

$$\langle x_n, t_n | x_{n-1}, t_{n-1} \rangle = \frac{1}{w(\Delta t)} \exp \left(\frac{iS(n, n-1)}{\hbar} \right) \quad (3)$$

$$\langle x_N, t_N | x_1, t_1 \rangle = \int_{x_1}^{x_N} \mathcal{D}(x(t)) \exp \left(i \int_{t_1}^{t_N} dt \frac{L_{\text{classical}}(x, \dot{x})}{\hbar} \right)$$

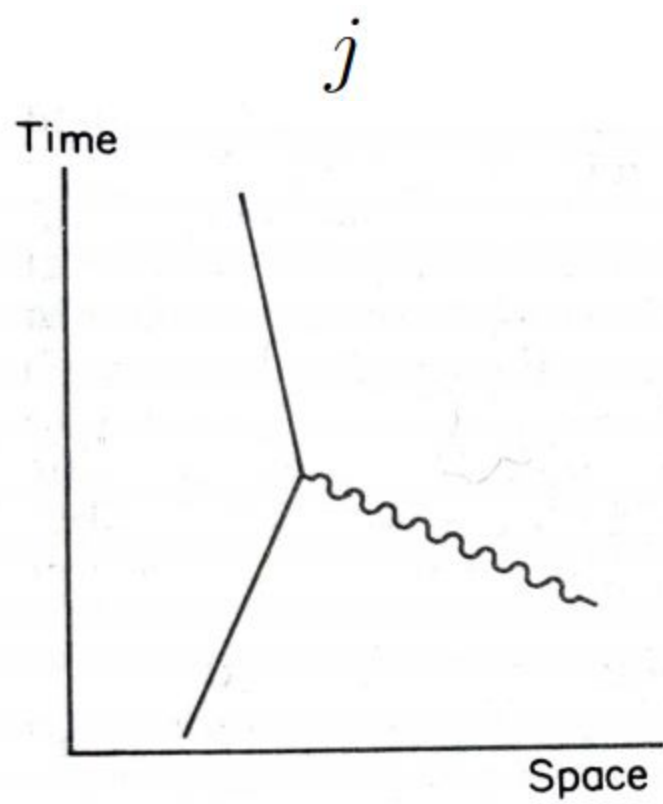
$$\int_{x_1}^{x_N} \mathcal{D}(x(t)) = \lim_{N \rightarrow \infty} \left(\frac{m}{2\pi i \hbar \Delta t} \right)^{\frac{N-1}{2}} \int dx_{N-1} \int dx_{N-2} \dots \int dx_2$$

If a photon exists at point A , given by the spacetime coordinates (x_1, y_1, z_1) at time t_1 , its probability of going to point B , given by the spacetime coordinates (x_2, y_2, z_2) at time t_2 , is inversely proportional to the spacetime interval⁷

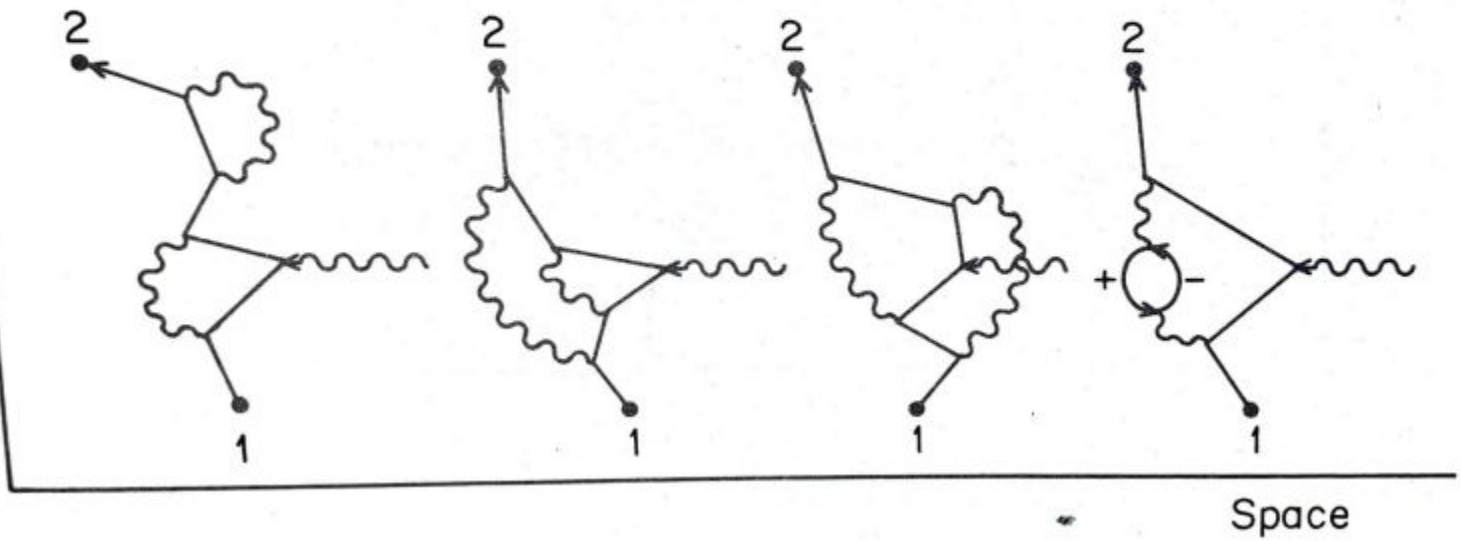
$$(x_2 - x_1)^2 + (y_2 - y_1)^2 + (z_2 - z_1)^2 - (t_2 - t_1)^2 \quad (8)$$

This probability is often denoted $P(A \text{ to } B)$ (“ P ” for ”photon”).

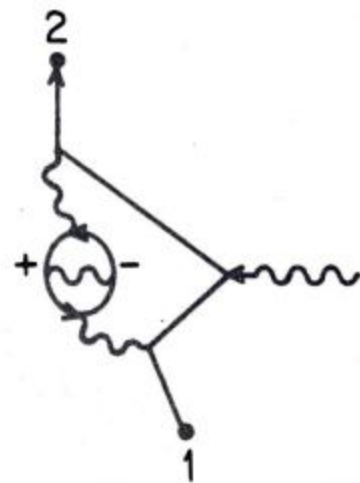
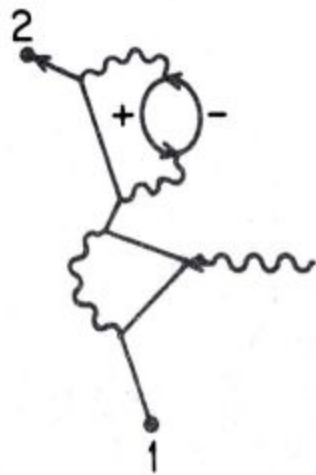
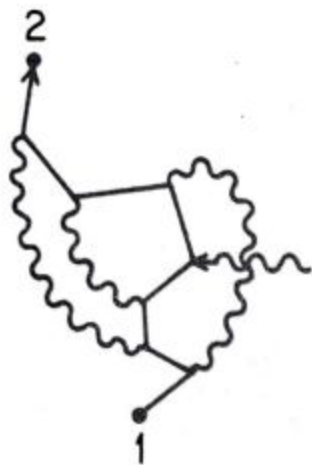
If an electron exists at point A , again given by the spacetime coordinates (x_1, y_1, z_1) at time t_1 , its probability of going to point B , again given by the spacetime coordinates (x_2, y_2, z_2) at time t_2 , is also related to the inverse of equation (8), *as well as* a number that we will call n . This is an empirically found number that helps us match the theory to experiment. This probability is often denoted $E(A \text{ to } B)$ (“ E ” for ”electrons”).



Time



Time



Space

1.00115965246

1.00115965221

Feynman notes: “If you were to measure the distance from Los Angeles to New York to this accuracy, it would be exact to the thickness of a human hair.”¹¹ Since Feynman’s day, this error has only decreased, making QED the most accurate physical theory.

References

David Griffiths. *Introduction to Quantum Mechanics*, 2nd ed. Pearson. 2015.

J.J. Sakurai. *Modern Quantum Mechanics*, 2nd ed. Cambridge University Press. 2017.

Richard P. Feynman. *QED: The Strange Theory of Light and Matter*. Princeton University Press. 1985.